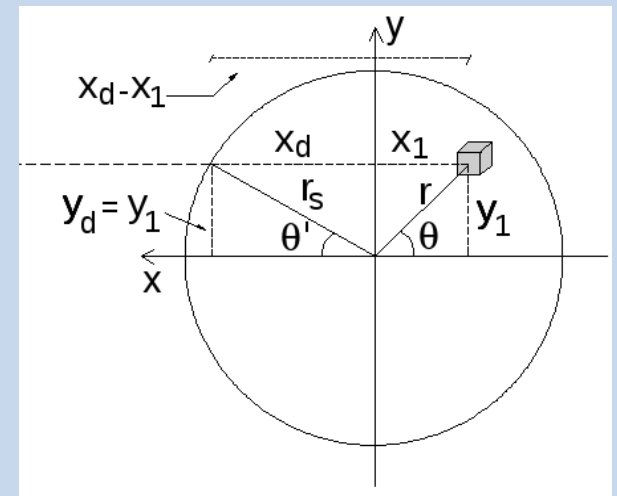


MODELLING THE PRODUCTION OF ELECTRONS AND X-RAYS FOR MONTE CARLO DOSIMETRY CALCULATIONS

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CONTEXT

PRIMARY AND SECONDARY EMISSION OF PARTICLES

⇒ **Environment (aerosols)**

⇒ **Accurate elemental characterization using PIXE analytical technique**

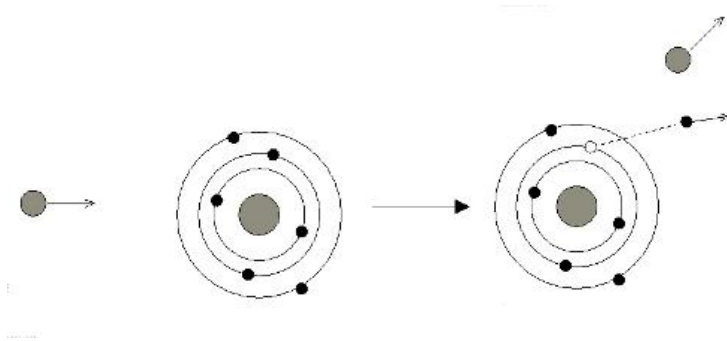
⇒ **Biomedical applications**

⇒ **Precise determination of the deposited energy in cells - Microdosimetry**
 (∞ number of emitted X-rays, Auger electrons, etc)

⇒ **Number of detected X-rays:**

$$N_{ij} = \frac{\Omega}{4\pi} \epsilon_j N T_j C_i \int_{E_{out}}^{E_{in}} \frac{\tau_j(x(E))}{S(E)} \sigma_{ij}^X(E) dE$$

ATOMIC EXCITATION - DEEXCITATION



⇒ H and He ion beams - PIXE (Particle induced X-ray Emission) analysis

⇒ β , α (He ions), γ particles, etc - nuclear decay

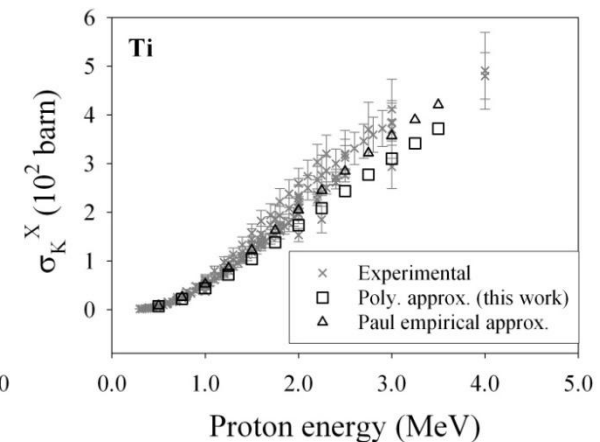
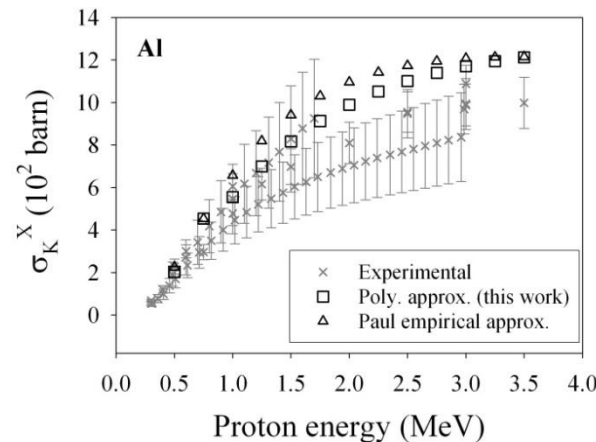
⇒ Atomic ionisation

⇒ Ionisation cross-sections

⇒ X-ray production cross-sections - ionisation cross-sections

⇒ Experimental values - large discrepancies between data sets

⇒ Theoretical values - experimental data dependency?



UNIVERSAL IONISATION CROSS-SECTIONS - H AND HE IONS

⇒ Theoretical approach - universal ionisation cross-sections

⇒ 7th order polynomial approximation

⇒ 0.1 MeV to 10 MeV (100 MeV)

$$\sigma_{n,o}^U(\xi_{n,o}^R, \zeta_{n,o}, \theta_{n,o}) = \frac{\sigma_{n,o}^{ECPSRR}(\xi_{n,o}^R, \zeta_{n,o}, \theta_{n,o})}{8\pi\alpha_0^2 Z_{proj}^2 C_{n,o}(x_q)} \eta_{n,o} Z_{n,o}^2 \theta_{n,o}^{1+c_{n,o}^U}$$

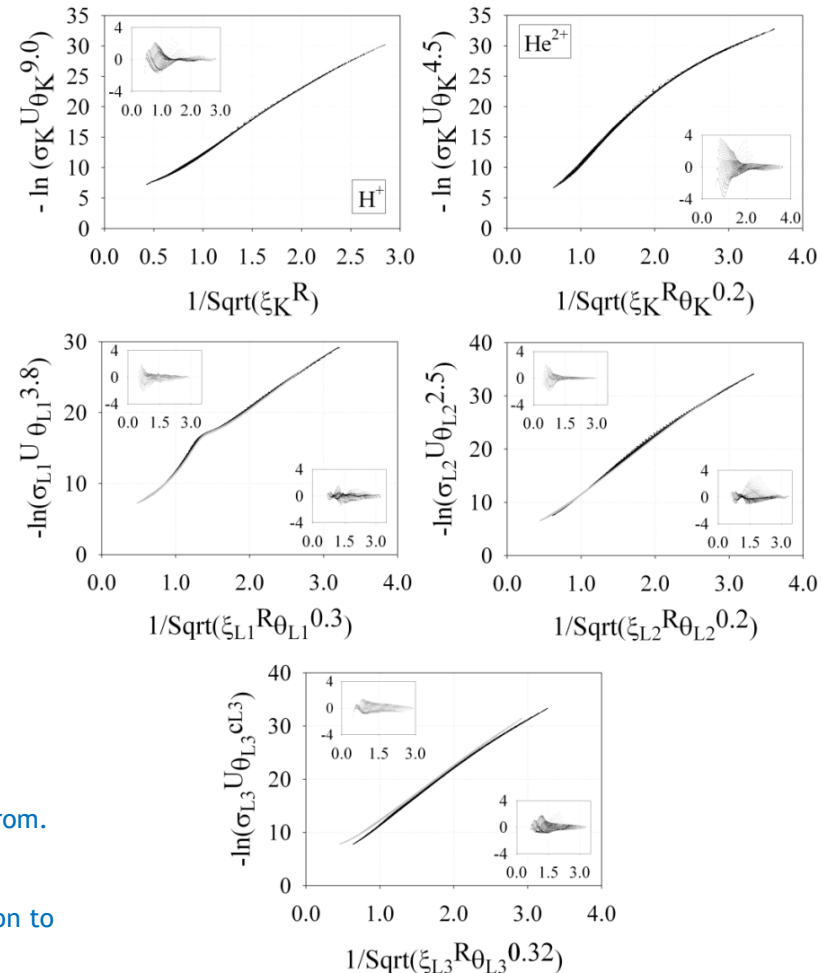
⇒ Allows exclusion of experimental data dependency

⇒ Partially implemented in Geant4 since 2012

References:

Taborda, A.; Chaves, P. C.; Reis, M. A., « Polynomial approximation to universal ionisation cross-sections of K and L shells induced by H and He ions », X-Ray Spectrom. 40:127, 2011

Taborda, A.; Chaves, P. C.; Carvalho, M. L.; Reis, M. A., « Polynomial approximation to universal M-shell ionisation cross-sections induced by H+ and He2+ ions », X-Ray Spectrom. 42:177, 2013



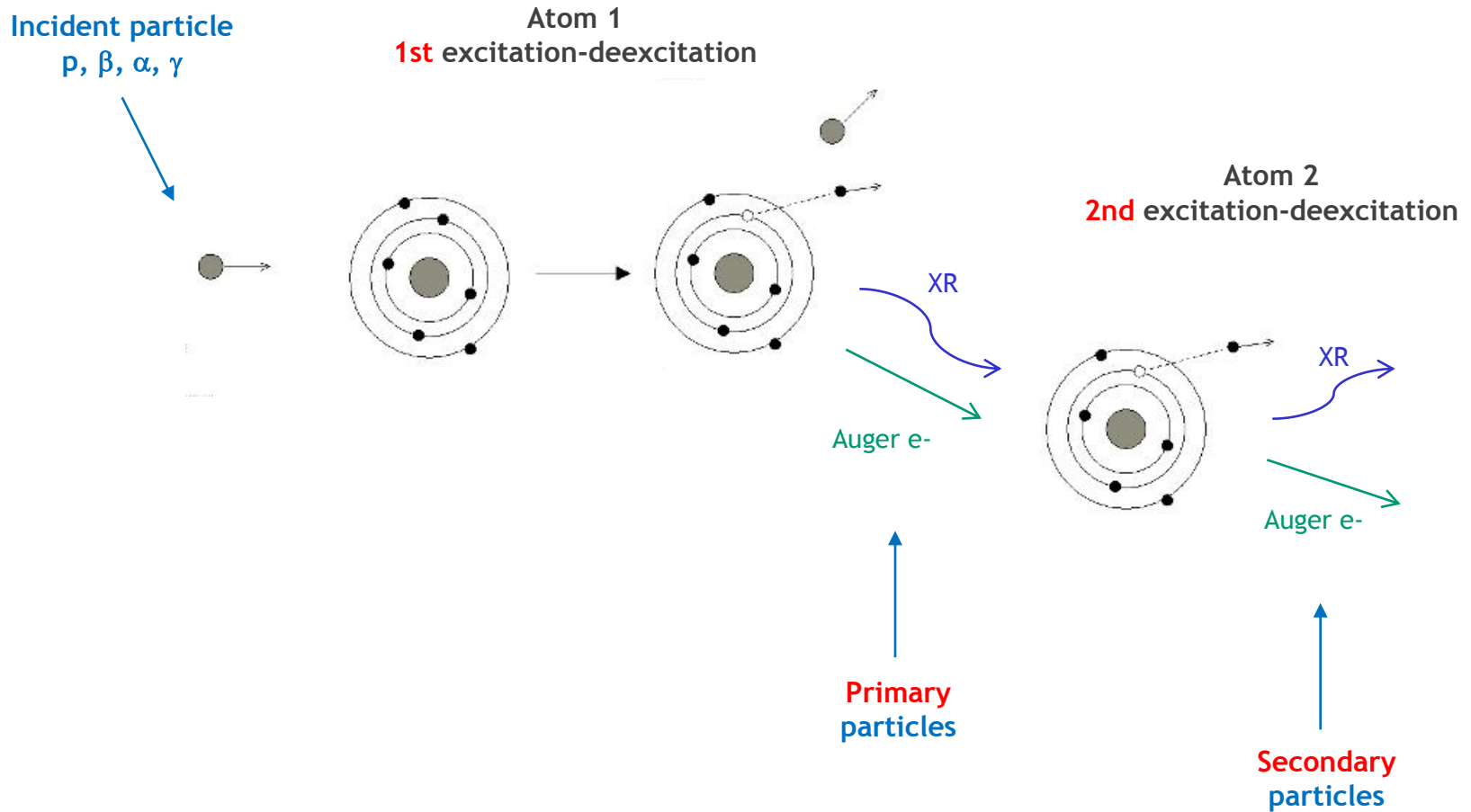
UNIVERSAL IONISATION CROSS-SECTIONS - HE IONS

- ⇒ Implemented in Geant4 → can be used in dosimetry calculations
- ⇒ Alpha particles emitting radionuclides, e.g.
- ⇒ Energies < 10 MeV

Radionuclide	α particle energy (MeV)
211At	5.8
213Bi	5.9
223Ra	5.7
225Ac	5.8
227Th	6.0
...	...

Energies values taken from RADAR - the radiation dose assessment resource

PRIMARY AND SECONDARY PARTICLES



SECONDARY X-RAYS

PENETRATION FUNCTION MODEL (REIS1996)

⇒ Number of X-rays B produced in a volume dV induced in dV by X-rays A

$$d\chi_{B,A} = P_A(x) R_{B,A} Q_A(r) T_B(x_1, r, \theta) dV$$

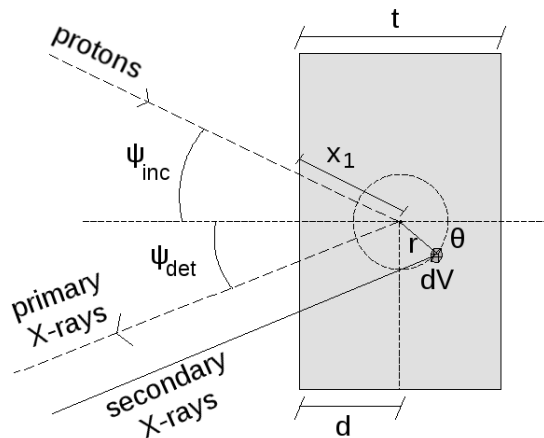
Probability of X-rays A produced by impact of particles at depth x

Probability that an absorbed X-ray A produces a secondary X-ray B

Density of probability of a X-ray A being absorbed at a distance r from the emission point

Transmission factor of the X-rays B from the volume dV up to the target surface, calculated in the detector direction

2D infinite sample



Secondary fluorescence emitted by the target is referred to the point where the emission of the primary radiation takes place

Reis, M. A., Alves, L. C., Jesus, A. P., Nucl. Instrum. Meth. B, 109/110: 134-138, 1996

SECONDARY X-RAYS

PENETRATION FUNCTION MODEL (REIS1996)

$$d\chi_{B,A} = P_A(x)R_{B,A}Q_A(r)T_B(x_1, r, \theta)dV$$

$$\Rightarrow \text{Probability of primary X-rays: } P_A(x) = \sigma_A^X(E_p),$$

$$\Rightarrow \text{Probability that } XR_A \text{ produces } XR_B: R_{B,A} = \frac{\sigma_B^{photo}(E_A)\omega_B k_B}{\mu(A)} f_B$$

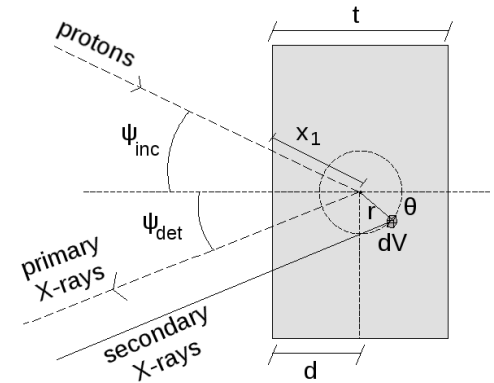
$$\Rightarrow \text{Absorption of } XR_A \text{ at } r \text{ from } x: Q_A(r) = \frac{\mu(A)}{4\pi r^2} e^{-\mu(A)r}$$

$$\Rightarrow \text{Transmission of } XR_B \text{ to surface: } T_B(r_1, r, \theta) = e^{-\frac{\mu(B)(x_1 \cos \Psi_{inc} + r \cos \theta)}{\cos \Psi_{det}}}$$

$$Q_{B,A}(x_1) = \int_{V_{target}} Q_A(r)T_B(x_1, r, \theta)dV$$

\Rightarrow Total number of secondary X-rays in thick sample:

$$\chi_{B,A} = P_A(x)R_{B,A}A(x_1) \left[I_1 + e^{\frac{\mu(B)}{\cos \Psi_{det}}} E_1[\mu(A)x \cos \Psi_{inc}] - E_1[g^+ x \cos \Psi_{inc}] \right]$$



GENERAL MODEL

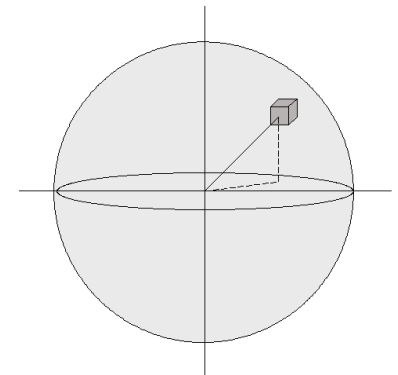
PENETRATION FUNCTION MODEL - advantages

⇒ Secondary fluorescence emitted by the target is referred to the point where the emission of the primary radiation takes place

⇒ For 2D infinite samples: analytical expressions

⇒ Especially useful for spherical samples

⇒ Number of particles B produced in a volume dV induced in dV by particles A (particles: X-rays, Auger e^-)



$$d\chi_{B,A} = P_A(x)R_{B,A}Q_A(r)T_B(x_1, r, \theta)dV$$

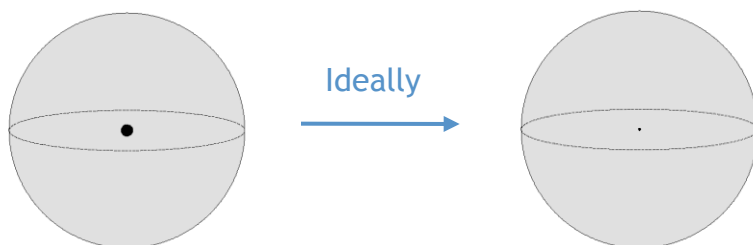
Determination of $P_A(x)$, R_{BA} , $Q_A(r)$ and $T_B(x_1, r, \theta)$!

GEOMETRY CONSIDERATIONS

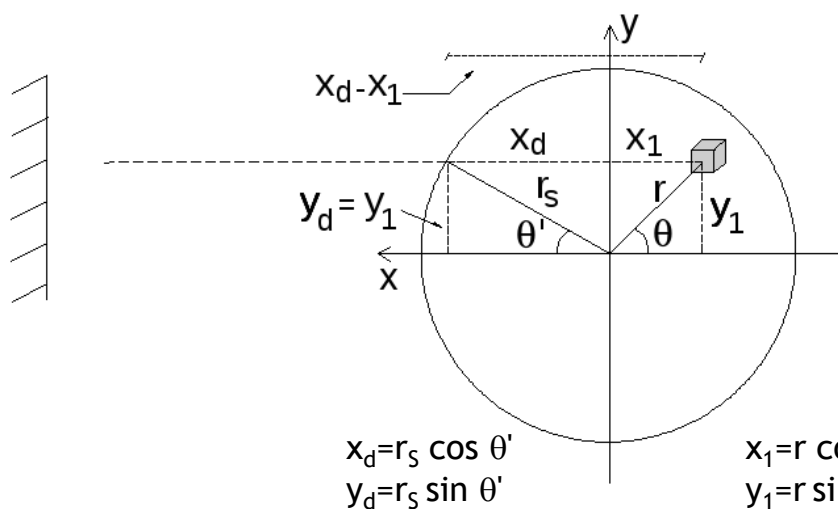
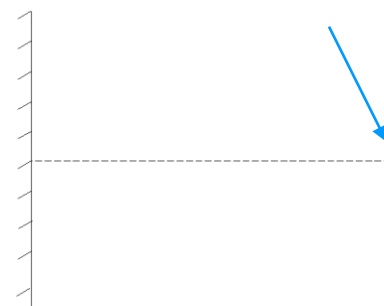
CENTRAL POINT EMISSION

⇒ Nanoparticle described by ideal geometrical model

Fe core: \varnothing 5 nm, Si coating: \varnothing 200 nm



Detector >> nanoparticle



Distance inside sphere: $x_d - x_1$

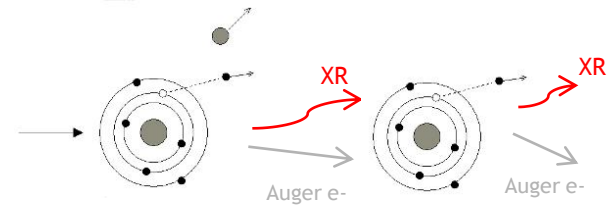
$y_d = y_1$

$\theta' = \arcsin(r \sin \theta / r_s)$

First step:

primary X-ray -> secondary X-ray
for homogeneous sphere

PRIMARY X-RAY



SECONDARY X-RAY

$$d\chi^{X_p, X_s} = P^{X_p}(E_H) R^{X_p, X_s} Q^{X_p}(r) T^{X_s}(r_S, r, \theta) dV$$

$$T^{X_s}(r_S, r, \theta) = e^{-\mu(E_{X_s})[x_d - x_1]} = e^{-\mu(E_{X_s})[\sqrt{r_S^2 - r^2 \sin^2 \theta} - r \cos \theta]}$$

$$Q^{X_p, X_s} = \int_{V_{\text{sphere}}} Q^{X_p}(r) T^{X_s}(r_S, r, \theta) dV$$

$$Q^{X_p, X_s} = \int_0^{r_S} \int_{-1}^1 \frac{\mu(E_{X_p})}{2} e^{-\mu(E_{X_p})r} e^{-\mu(E_{X_s})[\sqrt{r_S^2 - r^2(1-y^2)} + ry]} dy dr$$

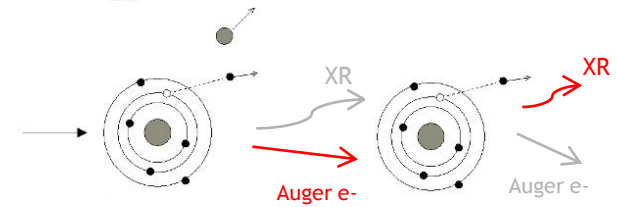
$$y = -\cos \theta$$

⇒ Number of secondary X-rays:

$$\begin{aligned} \chi^{X_p, X_s} &= \sigma^{X_p}(E_H) \frac{\sigma_s^{\text{photo}}(E_{X_p}) \omega_s k_{X_s}}{\mu(E_{X_p})} f_s \cdot \\ &\cdot \int_0^{r_S} \int_{-1}^1 \frac{\mu(E_{X_p})}{2} e^{-\mu(E_{X_p})r} e^{-\mu(E_{X_s})[\sqrt{r_S^2 - r^2(1-y^2)} + ry]} dy dr \end{aligned}$$

PRIMARY AUGER E⁻

SECONDARY X-RAY



⇒ Fluorescence yield and Auger yield

⇒ Fe(Z=26): $\omega_K=0.340$, $a_K=(1-\omega_K)=0.660$; Si(Z=14): $\omega_K=0.050$, $a_K=0.950$

⇒ Energy loss of primary particles

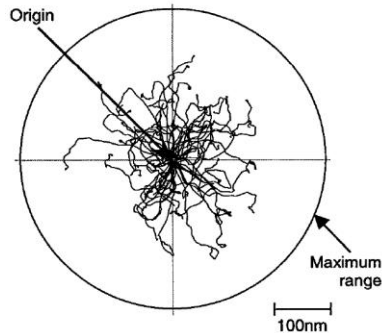


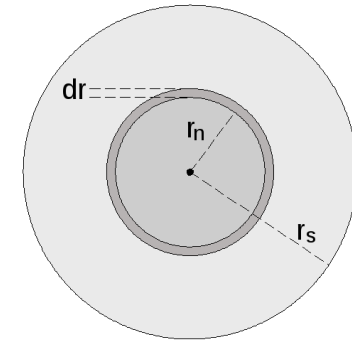
Figure: Electron trajectories in silicon at initial energy 3.5 keV (Fig 3 in Kataoka1999)

$$E(r) = E_0 - \int_0^{r_n} \frac{dT}{dr} dr$$

E_0 : initial energy of primary electrons

$-dT/dr$: total energy loss

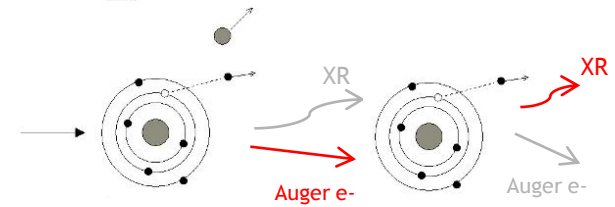
r_n : distance to dV



⇒ Number of electrons at position r:

$$N(r) = N_0 - \int_0^{r_n} N(r) \sigma^{I,e}(E_r) dr$$

PRIMARY AUGER E-



SECONDARY X-RAY

$$\Rightarrow d\chi^{e_p, X_s} = P^{e_p}(E_H, r) R^{e_p, X_s} Q^{e_p}(r) T^{X_s}(r_S, r, \theta) dV$$

\Rightarrow Absorption of primary electrons

$$\mathbf{I} \quad R^{e_p, X_s} \equiv 1$$

$$\mathbf{I} \quad Q^{e_p, X_s}(r) dV = \frac{\sigma_s^{I, e}(E_r) \omega_s k_{X_s} f_s}{4\pi r^2} dV$$

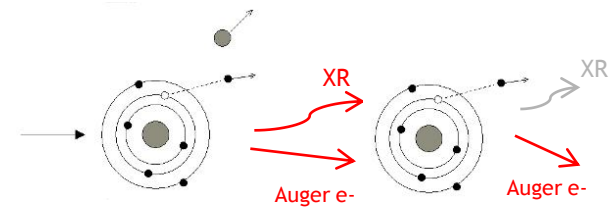
\Rightarrow Number of secondary X-rays:

$$\mathbf{I} \quad \chi^{e_p, X_s} = \int_{V_{\text{sphere}}} P^{e_p}(E_H, r) Q^{e_p, X_s}(r) T^{X_s}(r_S, r, \theta) dV$$

$$\mathbf{I} \quad \chi^{e_p, X_s} = \frac{\omega_s k_{X_s} f_s}{2} \int_0^{r_S} \int_{-1}^1 P^{e_p}(E_H, r) \sigma_s^{I, e}(E_r) e^{-\mu(E_{X_s}) [\sqrt{r_S^2 - r^2(1-y^2)} + ry]} dy dr$$

$$\mathbf{I} \quad P^{e_p}(E_H, r) = N(0) - \int_0^{r_n} N(r) \sigma_s^{I, e}(E_r) dr$$

SECONDARY AUGER E⁻



PRIMARY X-RAY OR PRIMARY AUGER E⁻

⇒ Transmission of secondary electrons

$$T^{e_s}(r_S, r, \theta) = e^{-\mu_m(E_r) \left[\sqrt{r_S^2 - r^2 \sin^2 \theta} - r \cos \theta \right]}$$

⇒ Primary X-ray - secondary Auger e⁻

$$\begin{aligned} \chi^{X_p, e_s} &= \sigma^{X_p}(E_H) \frac{\sigma_s^{photo}(E_{X_p})(1 - \omega_s)k_{e_s} f_s}{\mu(E_{X_p})} \cdot \\ &\cdot \int_0^{r_S} \int_{-1}^1 \frac{\mu(E_{X_p})}{2} e^{-\mu(E_{X_p})r} e^{-\mu_m(E_r) \left[\sqrt{r_S^2 - r^2(1-y^2)} + ry \right]} dy dr \end{aligned}$$

⇒ Primary Auger e⁻ - secondary Auger e⁻

$$\chi^{e_p, e_s} = \frac{(1 - \omega_s)k_{e_s} f_s}{2} \int_0^{r_S} \int_{-1}^1 P^{e_p}(E_H, r) \sigma_s^{I, e}(E_r) e^{-\mu_m(E_{X_s}) \left[\sqrt{r_S^2 - r^2(1-y^2)} + ry \right]} dy dr$$

APPLICATIONS

X-RAY EMITTING NANOPARTICLE

⇒ Fe core coated with Si

Central point: Fe(Z=26)
 Sphere: Si (Z=14)
 Sphere radius (r_s): 100 nm

$E_p=3.0$ MeV
 $E(\text{Si-KL23})=1.740$ keV

X-ray transition: KL23
 Auger transition: K-L23L23

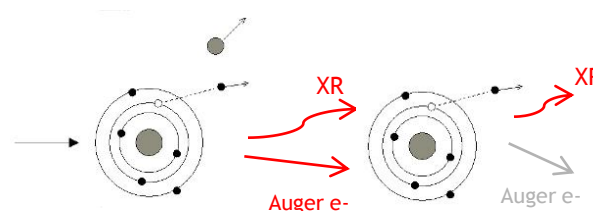
Fe X-ray production cross-section: 145.67 b
 Fe Auger e^- production cross-section: 213.99 b

Fe fluorescence yield: 0.34
 Si fluorescence yield: 0.05

⇒ How many Si KL23 X-rays exit the nanoparticle?

$$\chi_{Si,Fe}^{X,X} = 0.0097$$

$$\chi_{Si,Fe}^{e,X} = 0.0140$$



M. O. Krause, J. Phys. Chem. Ref. Data, 8: 307-327, 1979
 J. H. Scofield, At. Data Nucl. Data Tables, 14: 121-137, 1974
 F. P. Larkins, At. Data Nucl. Data Tables, 20: 311-387, 1977
 D. L. Walters, C. P. Bhalla, Atomic Data, 3: 301-315, 1971

APPLICATIONS

AUGER E⁻ EMITTING NANOPARTICLES

⇒ Fe core coated with Si

Central point: Fe(Z=26)
 Sphere: Si (Z=14)
 Sphere radius (r_s): 100 nm

$E_p=3.0$ MeV
 $E(\text{Si-K-L23KL23})=1.846$ keV

X-ray transition: KL23
 Auger transition: K-L23L23

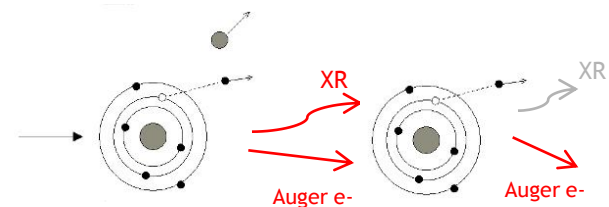
Fe X-ray production cross-section: 145.67 b
 Fe Auger e⁻ production cross-section: 213.99 b

Fe fluorescence yield: 0.34
 Si fluorescence yield: 0.05

⇒ How many Si K-L23L23 Auger electrons exit the nanoparticle?

$$\chi_{Si,Fe}^{X,e} = 0.1239$$

$$\chi_{Si,Fe}^{e,e} = 0.1778$$



M. O. Krause, J. Phys. Chem. Ref. Data, 8: 307-327, 1979
 J. H. Scofield, At. Data Nucl. Data Tables, 14: 121-137, 1974
 F. P. Larkins, At. Data Nucl. Data Tables, 20: 311-387, 1977
 D. L. Walters, C. P. Bhalla, Atomic Data, 3: 301-315, 1971

SUMMARY AND CONCLUSIONS

- ⇒ Secondary X-ray and electron emission expressions have been obtained for primary X-rays and Auger electrons for an ideal spherical model
- ⇒ Can be used for PIXE analysis (impact of particles inducing the primary particles emission) or consider an emitting radionuclide
- ⇒ As a simple model it can be used for the validation and intercomparison with results from Monte Carlo codes
- ⇒ Number of electrons or X-rays leaving the nanoparticle fundamental for the accurate account of deposited energy (and possible damage to cells)
- ⇒ To be further developed to become a more realistic model: one direction for exiting the nanoparticle and only one step of electron cascade is not realistic

THANK YOU

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